SOME REMARKS ON THE PAPER, ENTITLED "FRACTIONAL AND OPERATIONAL CALCULUS WITH GENERALIZED FRACTIONAL DERIVATIVE OPERATORS AND MITTAG-LEFFLER TYPE FUNCTIONS" BY Z. TOMOVSKI, R. HILFER AND H. M. SRIVASTAVA

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ABSTRACT. In this note, we would like to bring the readers' attention toward the fact that [9, Section 3, Eq.(3.2)] is not a solution of the fractional differential equation [9, Section 3, Eq.(3.1)].

Keywords: generalized Mittag-Leffer type functions, fractional differential equations, Dirac delta function.

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As is well known (see, for example, [1-5], [8], a great deal of attention has recently been paid to the solution of differential equations involving both ordinary and partial derivatives of fractional order. However, we propose to show here that [9, Section 3, Eq.(3.2)] is not a solution of the fractional differential equation [9, Section 3, Eq.(3.1)].

We begin by recalling the following formulas:

$$\left(D_{0+}^{\alpha}f\right)(x) = \lambda f(x) \tag{1}$$

and

$$f(x) = x^{1-\alpha} E_{\alpha,\alpha} \left(\lambda x^{\alpha}\right).$$
⁽²⁾

Here the equations (1) and (2) correspond to Eq. (3.1) and Eq. (3.2) in [9]. The Mittag-Lefler type function occurring in the equation (2) is defined by

$$E_{\mu,\nu}(z) := \sum_{n=0}^{\infty} \frac{z^n}{\Gamma(\mu n + \nu)}.$$
(3)

Making use of Eq. (2) and the definition (3), we observe that

$$\begin{split} f\left(x\right) &= x^{1-\alpha}\sum_{n=0}^{\infty}\frac{(\lambda x^{\alpha})^{n}}{\Gamma\left(\alpha n+\alpha\right)} = \sum_{n=0}^{\infty}\frac{\lambda^{n}\ x^{\alpha n+1-\alpha}}{\Gamma\left(\alpha n+\alpha\right)} \\ &= \sum_{n=0}^{\infty}\frac{\lambda^{n}}{\Gamma\left(\alpha n+\alpha\right)}\ \cdot\ \Gamma\left(\alpha n+2-\alpha\right)\ \cdot\ \frac{x^{\alpha n+1-\alpha}}{\Gamma\left(\alpha n+2-\alpha\right)} \end{split}$$

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so that

$$\left(D_{0+}^{\alpha} f \right)(x) = \sum_{n=0}^{\infty} \frac{\lambda^n \, \Gamma\left(\alpha n + 2 - \alpha\right)}{\Gamma\left(\alpha n + \alpha\right)} \, D_{0+}^{\alpha} \left(\frac{x^{\alpha n + 1 - \alpha}}{\Gamma\left(\alpha n + 2 - \alpha\right)} \right)$$
$$= \sum_{n=0}^{\infty} \frac{\lambda^n \, \Gamma\left(\alpha n + 2 - \alpha\right)}{\Gamma\left(\alpha n + \alpha\right)} \, \frac{x^{\alpha n + 1 - 2\alpha}}{\Gamma\left(\alpha n + 2 - 2\alpha\right)} \neq \lambda f(x) \, .$$

 $f(x) = x^{\alpha - 1} E_{\alpha, \alpha} \left(\lambda \, x^{\alpha} \right),$

Hence, clearly, (2) is *not* a solution of (1).

If we take

then

$$f(x) = \sum_{n=0}^{\infty} \frac{\lambda^n x^{\alpha n + \alpha - 1}}{\Gamma(\alpha n + \alpha)}.$$

In this case, we readily see that

$$\left(D_{0+}^{\alpha} f \right)(x) = D_{0+}^{\alpha} \left(\frac{x^{\alpha n + \alpha - 1}}{\Gamma(\alpha n)} \right) \Big|_{n=0} + \sum_{n=1}^{\infty} \frac{\lambda^n}{\Gamma(\alpha n + \alpha)} D_{0+}^{\alpha} \left(\frac{x^{\alpha n + \alpha - 1}}{\Gamma(\alpha n + \alpha)} \right)$$

$$= \frac{x^{-1}}{\Gamma(0)} + \sum_{n=1}^{\infty} \frac{\lambda^n x^{\alpha n - 1}}{\Gamma(\alpha n)},$$

$$(5)$$

where, obviously, [7]

$$\frac{x^{-1}}{\Gamma(0)} = \frac{x^{-1}}{(-1)!} = 0 \qquad \text{whenever} \qquad x \neq 0.$$
 (6)

Since the Riemann-Liouville fractional calculus is based upon a definite integral which is taken over a *non-empty* interval (0, x), we can tacitly assume that x > 0. Thus, for x > 0, we find from (5) and (6) that

$$(D_{0+}^{\alpha} f)(x) = \sum_{n=1}^{\infty} \frac{\lambda^n x^{\alpha n-1}}{\Gamma(\alpha n-1)}$$

=
$$\sum_{n=0}^{\infty} \frac{\lambda^{n+1} x^{\alpha(n+1)-1}}{\Gamma(\alpha(n+1))}$$

=
$$\lambda x^{\alpha-1} E_{\alpha,\alpha} (\lambda x^{\alpha}) = \lambda f(x) \qquad (x > 0).$$

It follows that the *correct* solution of (1) is given by (4) under the *explicitly-stated* condition that x > 0.

The above-mentioned error was reproduced in a relatively more recent survey-cum-expository article [8] on the theory and applications of the Mittag-Leffler type functions which are associated with various operators of fractional calculus.

Finally, we turn to the familiar fact that the impulse or distributional (or generalized) function $\delta(x)$, which is popularly known as the *Dirac delta function*, is traditionally defined, for any suitably-constrained continuous function $\varphi(x)$, by (see, for details, [2] and [6]; see also [5] and [7]):

(4)

$$\delta(x) = 0$$
 $(x \neq 0)$ and $\int_{-\infty}^{\infty} \delta(x)\varphi(x) \, \mathrm{d}x = \varphi(0),$ (7)

so that, in particular, we have

$$\delta(x) = 0$$
 $(x \neq 0)$ and $\int_{-\infty}^{\infty} \delta(x) \, \mathrm{d}x = 1.$ (8)

In light of Eq. (6), it is possible to set [7]

$$\delta(x) = \frac{x^{-1}}{\Gamma(0)} = \frac{x^{-1}}{(-1)!} = 0 \quad \text{whenever} \quad x \neq 0.$$
(9)

However, for obvious reasons, Eq. (9) cannot be construed to define the Dirac delta function $\delta(x)$, simply because it does not satisfy the necessary *second* requirement in the definition (8). Consequently, instead of the Dirac delta function $\delta(x)$, use should be (and has been) made in Eq. (5) of the quotient involved in Eq. (6) and Eq. (9).

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